Bell's local causality in local classical and quantum theory

Gábor Hofer-Szabó Research Centre for the Humanities, Budapest

Péter Vecsernyés Wigner Research Centre for Physics, Budapest

Main message

- Bell's notion of local causality presupposes a framework integrating probabilistic and spatiotemporal entities.
- Our aim is to develop such a framework called local physical theory.

- I. What is a local physical theory?
- **II.** Bell's local causality in a local physical theory
- **III.** Common Cause Principle
- IV. Bell inequalities

Discretized two dimensional Minkowski spacetime:



- \mathcal{M} : globally hyperbolic spacetime with symmetries \mathcal{P}
- ${\cal K}$: covering collection of bounded, globally hyperbolic regions of ${\cal M}$
- (\mathcal{K}, \subseteq) : directed poset
- $\mathcal{P}_{\mathcal{K}}$: subgroup of \mathcal{P} leaving \mathcal{K} invariant

- Definition. A P_K-covariant local physical theory (LPT) is a net K ∋ V → N(V) associating von Neumann algebras to spacetime regions which satisfies
 - 1. isotony,
 - 2. microcausality,
 - 3. covariance.

• Isotony: if $V_1 \subset V_2$, then $\mathcal{N}(V_1)$ is a unital subalgebra of $\mathcal{N}(V_2)$



• Microcausality (Einstein causality): $[\mathcal{N}(V_A), \mathcal{N}(V_B)] = 0$



• Covariance: covariant group homomorphism on the net $\alpha_g(\mathcal{N}(V)) = \mathcal{N}(g \cdot V)$



Remarks:

- C^* -algebras, von Neumann algebras, σ -algebras
- Quasilocal algebra A: the inductive limit C*-algebra of the net
- A is commutative: **local classical theory (LCT)**
- A is noncommutative: local quantum theory (LQT)

• Causal complement: V'



• Domain of dependence: V''



Examples:

- 1. Deterministic LCT
- 2. Stochastic LCT
- 3. Deterministic LQT
- 4. Stochastic LQT (not known)

1. Deterministic LCT



Local algebras:



Local algebras:





















2. Stochastic LCT



Stochastic dynamics: with probability *p*



• Stochastic dynamics: with probability 1 - p



Stochastic dynamics:



Stochastic dynamics: with probability *p*



• Stochastic dynamics: with probability 1 - p



3. **Deterministic LQT** by imposing anticommutation relation between neighbours



• "A theory will be said to be locally causal if the probabilities attached to values of local beables in a space-time region V_A are unaltered by specification of values of local beables in a space-like separated region V_B , when what happens in the backward light cone of V_A is already sufficiently specified, for example by a full specification of local beables in a space-time region V_C ." (Bell, 1990/2004, p. 239-240)



Basic terms:

- 1. "The **beables** of the theory are those entities in it which are, at least tentatively, to be taken seriously, as corresponding to something real."
- 2. "there *are* things which **do go faster than light**. British sovereignty is the classical example. When the Queen dies in London (long may it be delayed) the Prince of Wales, lecturing on modern architecture in Australia, becomes instantaneously King."
- 3. "*Local* beables are those which are definitely associated with particular space-time regions. The electric and magnetic fields of classical electromagnetism, $\mathbf{E}(t, x)$ and $\mathbf{B}(t, x)$ are again examples."

Basic terms:

4. "It is important that region V_C completely shields off from V_A the overlap of the backward light cones of V_A and V_B ."



Basic terms:

5. "And it is important that events in V_C be **specified completely**. Otherwise the traces in region V_B of causes of events in V_A could well supplement whatever else was being used for calculating probabilities about V_A ."

Translation:

- "local beable" \longrightarrow element of a local von Neumann algebra
- "complete specification" \longrightarrow an atomic element of a local von Neumann algebra

• "completely shielder-off region":

(i) $V_C \subset J_-(V_A)$ (ii) $V_A \subset V_C''$ (iii) $V_C \subset V_B''$



- Definition. A LPT is called (Bell) locally causal, if
 - for any pair of projections $A \in \mathcal{N}(V_A)$ and $B \in \mathcal{N}(V_B)$ supported in spacelike separated regions, and
 - for every locally normal and faithful state ϕ establishing a correlation between A and B, $\phi(AB) \neq \phi(A)\phi(B)$, and
 - for any spacetime region V_C satisfying Requirements
 (i)-(iii), and
 - for any *atomic event* C_k in $\mathcal{N}(V_C)$:

 $\frac{\phi(C_k A B C_k)}{\phi(C_k)} = \frac{\phi(C_k A C_k)}{\phi(C_k)} \frac{\phi(C_k B C_k)}{\phi(C_k)}$

Question:

When is a LPT locally causal?

• Local primitive causality: $\mathcal{N}(V) = \mathcal{N}(V'')$ for any $V \in \mathcal{K}$



Local primitive causality: holds in deterministic LCTs



Local primitive causality: does not hold in stochastic
 LCTs



Proposition:

 Any atomic LPT satisfying local primitive causality is locally causal.

But...

how can a LQT be locally causal if local causality implies the Bell inequalities which are violated for certain quantum correlations?

Reichenbach's Common Cause Principle (CCP): If there is a correlation between two events and there is no direct causal (or logical) connection between them, then there always exists a common cause of the correlation.

- Correlation: $\phi(AB) \neq \phi(A)\phi(B)$
- Common cause: partition $\{C_k\}_{k \in K}$ of the unit

$$\frac{\phi(C_k A B C_k)}{\phi(C_k)} = \frac{\phi(C_k A C_k)}{\phi(C_k)} \frac{\phi(C_k B C_k)}{\phi(C_k)}$$

- Correlation: $\phi(AB) \neq \phi(A)\phi(B)$
- Common cause: partition $\{C_k\}_{k \in K}$ of the unit

$$\frac{\phi(C_k A B C_k)}{\phi(C_k)} = \frac{\phi(C_k A C_k)}{\phi(C_k)} \frac{\phi(C_k B C_k)}{\phi(C_k)}$$

- Commuting / Noncommuting common cause: $\{C_k\}_{k\in K}$ is commuting / not commuting with A and B
- Nontrivial common cause:

 $C_k \not\leq A, A^{\perp}, B \text{ or } B^{\perp} \text{ for some } k \in K$

Common Cause Principle:



Weak Common Cause Principle:



Local causality:



What are the similarities and the differences between local causality and the Common Cause Principle?

Similarities:

- 1. Both local causality and the CCPs are properties of a LPT represented by a net $\{\mathcal{N}(V), V \in \mathcal{K}\}$.
- 2. The core mathematical requirement of both principles is the screening-off condition:

$$\frac{\phi(C_k A B C_k)}{\phi(C_k)} = \frac{\phi(C_k A C_k)}{\phi(C_k)} \frac{\phi(C_k B C_k)}{\phi(C_k)}$$

3. The **Bell inequalities** can be derived from both principles. (But see below.)

Differences:

- 1. For local causality the screening-off condition is required for **every** atomic event. For the CCPs it is required only for events of **one** partition.
- 2. For local causality the screening-off condition is required only for **atomic** events. For the CCPs one is looking for **nontrivial** common causes.
- For local causality screener-offs are localized
 'asymmetrically' in the past of V_A (or V_B). For the CCP they are localized 'symmetrically' in the joint / common past of V_A and V_B.

• A nice **parallelism**:

 $\begin{array}{rcl} \mbox{Local causality} & \Longrightarrow & \mbox{Bell inequalities} \\ \mbox{Common Cause Principle} & \Longrightarrow & \mbox{Bell inequalities} \end{array}$

- Set of correlations: $\phi(A_m B_n) \neq \phi(A_m) \phi(B_n)$
- Joint common cause: partition $\{C_k\}_{k \in K}$ of the unit

$$\frac{\phi(C_k A_m B_n C_k)}{\phi(C_k)} = \frac{\phi(C_k A_m C_k)}{\phi(C_k)} \frac{\phi(C_k B_n C_k)}{\phi(C_k)}$$

• Reduced state: $\phi_{\{C_k\}}(X) := \sum_k \phi(C_k X C_k)$

Proposition:

• Joint common cause \implies Bell inequalities for the reduced state $\phi_{\{C_k\}}$

 $-1 \leqslant \phi_{\{C_k\}}(A_1B_1 + A_1B_2 + A_2B_1 - A_2B_2 - A_1 - B_1) \leqslant 0$

• Joint common cause + commutativity \implies Bell inequalities for the original state ϕ

 $-1 \leqslant \phi(A_1B_1 + A_1B_2 + A_2B_1 - A_2B_2 - A_1 - B_1) \leqslant 0$

Proposition:

• Local causality \implies Bell inequalities for the reduced state $\phi_{\{C_k\}}$

 $-1 \leqslant \phi_{\{C_k\}}(A_1B_1 + A_1B_2 + A_2B_1 - A_2B_2 - A_1 - B_1) \leqslant 0$

• Local causality + commutativity \implies Bell inequalities for the original state ϕ

 $-1 \leqslant \phi(A_1B_1 + A_1B_2 + A_2B_1 - A_2B_2 - A_1 - B_1) \leqslant 0$

Conclusions

- Bell's notion of local causality presupposes a clear-cut framework integrating probabilistic and spatiotemporal entities. This goal can be met by introducing the notion of a LPT.
- In this general framework one can define Bell's notion of local causality and show sufficient conditions on which a LPT will be locally causal.
- There is a nice parallelism between local causality and the CCPs: Bell inequalities cannot be derived from neither unless the LPT is classical or the common cause is commuting.

References

- J.S. Bell, Speakable and Unspeakable in Quantum Mechanics, Cambridge: Cambridge University Press, (2004).
- J. Butterfield, "Stochastic Einstein Locality Revisited," Brit. J. Phil. Sci., 58, 805-867, (2007).
- J. Earman and G. Valente, "Relativistic causality in algebraic quantum field theory," (manuscript), (2014).
- R. Haag, Local Quantum Physics, (Springer Verlag, Berlin, 1992).
- J. Henson, "Non-separability does not relieve the problem of Bell's theorem," Found. Phys., **43**, 1008-1038 (2013).
- G. Hofer-Szabó, M. Rédei and L. E. Szabó, *The Principle of the Common Cause*, Cambridge: Cambridge University Press, 2013
- G. Hofer-Szabó and P. Vecsernyés, "Bell inequality and common causal explanation in algebraic quantum field theory," Stud. Hist. Phil. Mod. Phys., 44 (4), 404-416 (2013b).
- G. Hofer-Szabó and P. Vecsernyés, "On Bell's local causality in local classical and quantum theory," (in preparation) (2014).
- T. Norsen, "J.S. Bell's concept of local causality," Am. J. Phys, 79, 12, (2011).
- M. Rédei, "Reichenbach's Common Cause Principle and quantum field theory," Found. Phys., 27, 1309-1321 (1997).
- M. Rédei and J. S. Summers, "Local primitive causality and the Common Cause Principle in quantum field theory," Found. Phys., **32**, 335-355 (2002).